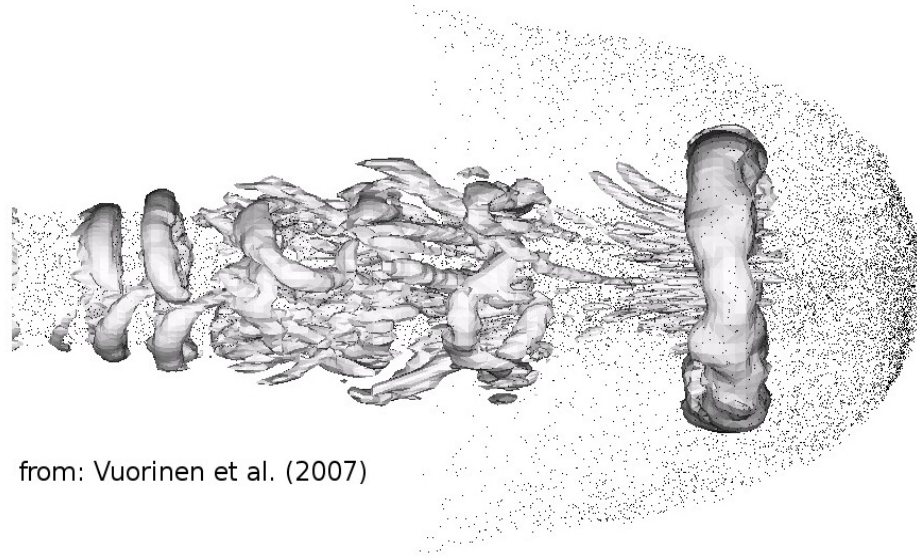


# Large-Eddy Simulation of Particle Size Distribution Effects on Turbulence in Sprays – Some Fundamental Phenomena



from: Vuorinen et al. (2007)

Spray Workshop, April 12th, 2008  
Cobo Hall, Detroit

Ville Vuorinen, ICEL/TKK, Finland  
Martti Larmi, ICEL/TKK, Finland  
Laszlo Fuchs, KTH, Sweden



# About This Presentation

- The presentation summarizes part of the work done by the authors during the year 2007
- Based on the papers
  - Vuorinen, Larmi, Fuchs, Large-Eddy Simulation of Spray-Originated Turbulence Production and Dissipation, ICMF-2007, Leipzig, (2007).
  - Vuorinen, Larmi, Fuchs, Large-Eddy Simulation of Droplet Size Distribution Effects on Turbulence in Sprays, AIAA-2008, Grand Sierra Resort, Reno (2008).

# Contents

- Background on Particles in Multiphase Flow and Objectives
- Assumptions on the Particulate Phase
- Problem Setup
- Large-Eddy Simulation (LES)
- Results
- Conclusions

# Background: Particles in Turbulence

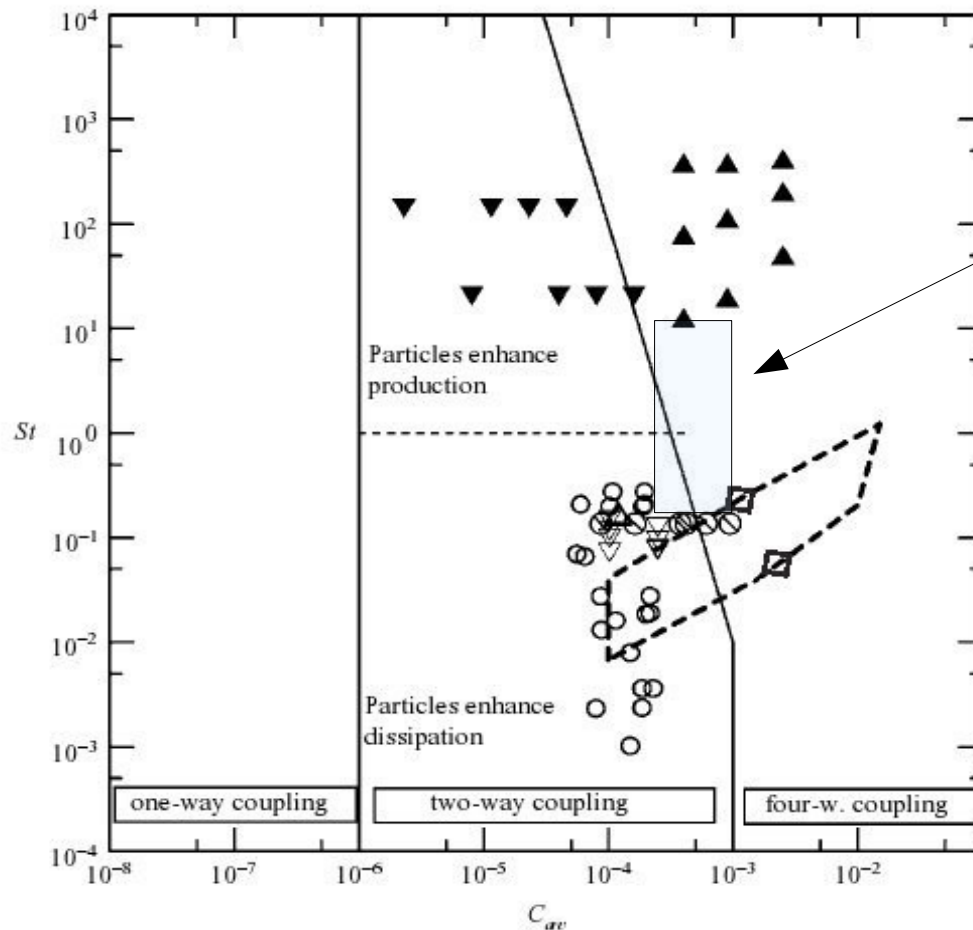
- the characterization by Elgobashi (1994) on 'generic' particle laden flows: particle size and volume fraction determine turbulence dissipation/production.

Particle Stokes number:

$$St_p = \tau_p \frac{U}{\delta}$$

Particle Momentum relaxation time:

$$\tau_p = \frac{\rho_p d^2}{18 \rho_g \nu}$$



this study

size

volume fraction

These definitions imply that small particles may follow a range of time frequencies.

# Objectives

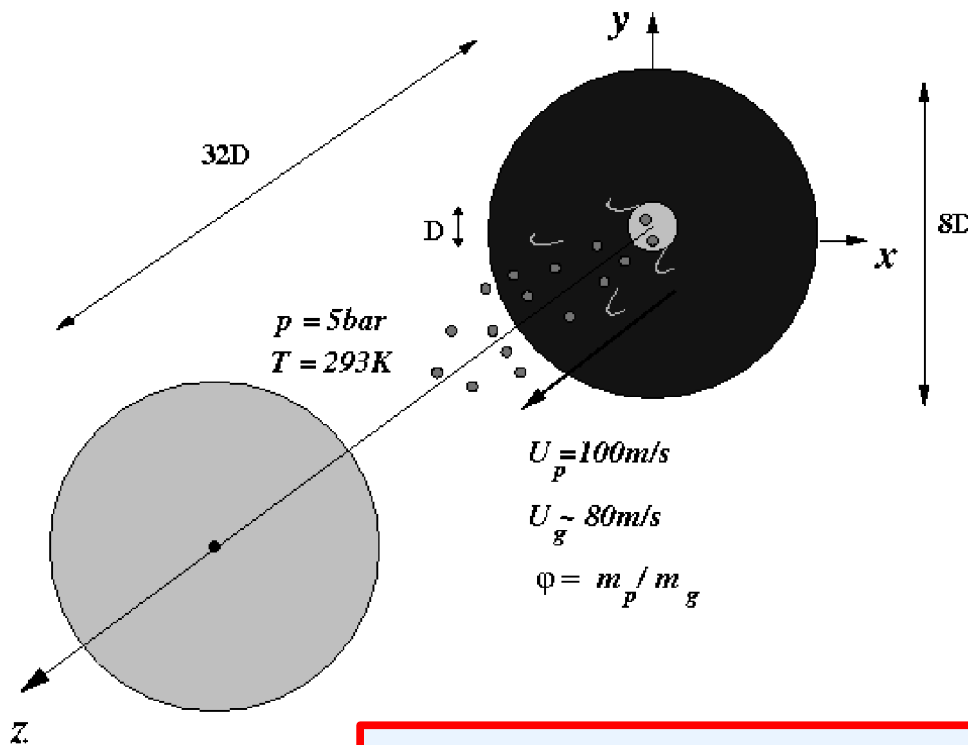
- Discuss the role of droplet size in forming the spray dynamics
- Discuss the connection between small scale interactions and large scale observations and point out that the PDF's of droplet velocity and slip velocity explain the spray behavior
- Look at flow structures and preferential concentration
- Discuss the characteristics of turbulent diffusivity in particle laden flows

# Assumptions on Particulate Phase

- the Lagrangian Particle Tracking approach (**LPT**)
- particles do not displace fluid
- particles are spherical and do not break
- particles do not interact (two way coupling)
- the stochastic “parcel” concept
- parcels couple to flow field via slip velocity
- gaseous phase is affected by the forces the parcels exert in each computational cell
- the gas velocity “seen” by parcel computed by linear interpolation to the parcel position

# Problem Setup

- Particles enter a gas jet
- slip velocity  $+0.25U_{\text{exit}}$
- two mass loadings (i.e. 0.3 and 0.6) studied corresponding to strong and very strong two way coupling (i.e. 0.001 and 0.002 volume fractions)
- particles are distributed according to a size distribution
- injection time = 1.5ms
- inlet diameter  $D=2\text{mm}$



$$\varphi = \frac{m_f}{m_g} = \frac{\text{injected spray mass}}{\text{injected gas mass}}$$

mass loading ratio

# Large-Eddy Simulation (LES)

- the full compressible Navier-Stokes is solved by a numerical algorithm (2<sup>nd</sup> order accurate in space, 1<sup>st</sup> order time)

$$\frac{\partial \rho u_i}{\partial t} + \frac{\partial (\rho u_i u_j)}{\partial x_j} = - \frac{\partial}{\partial x_j} (-p \delta_{ij} + \sigma_{ij}) + \mathcal{M}_{spray}$$

$$\mathcal{M}_{spray} = -C_D (u_i - u_{d,i}) |u_i - u_{d,i}| \delta_{r,r_d}$$

- “implicit LES”: filtering done by the discretization scheme
- an additional transport equation is solved for a passive scalar
- gaseous phase receives momentum from the particle phase via **the slip velocity**  $\delta W = u_{particle} - U_{gas}$

# Parallel Simulations

- simulations carried out with OpenFOAM open source control volume code
- mesh contains 3.5M cells – cell resolution as small as 40um in the shear layer and center of the jet
- decomposition onto 32 processors
- one simulation takes about 4 days
- CFL < 0.12
- further details in the referred manuscripts

# Simulated Cases

Stokes number as referred to SMD

Case	SMD	SMD/D	$\varphi$	distribution	$V_p/V_g$	$St_{SMD}$	$N_{p,SMD}$
A1	2.0 $\mu m$	0.001	0.3	monodisp.	0.001	0.1	125
A2	3.5 $\mu m$	0.00175	0.3	polydisp.	0.001	0.25	23
A3	7.0 $\mu m$	0.0035	0.3	polydisp.	0.001	1	3
A4	10.0 $\mu m$	0.005	0.3	polydisp.	0.001	2	1
B1	2.0 $\mu m$	0.001	0.6	monodisp.	0.002	0.1	250
B2	3.5 $\mu m$	0.00175	0.6	polydisp.	0.002	0.25	46
B3	7.0 $\mu m$	0.0035	0.6	polydisp.	0.002	1	6
B4	10.0 $\mu m$	0.005	0.6	polydisp.	0.002	2	2

mass loading 0.3

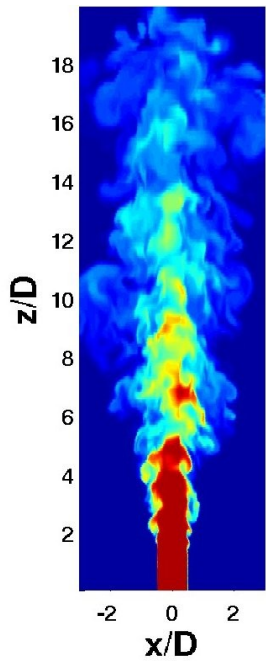
mass loading 0.6

4 different particle mean diameters

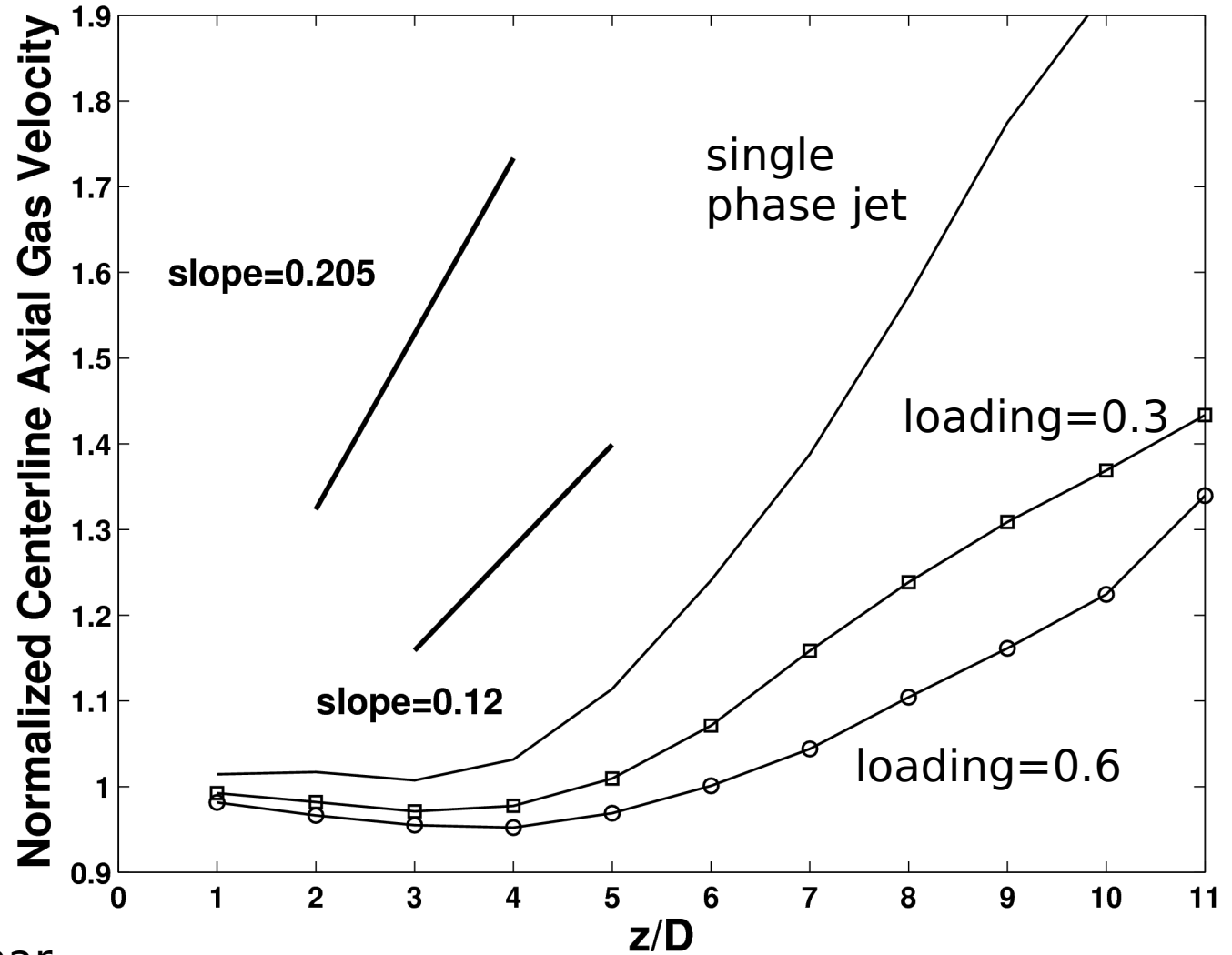
monodisperse = single size drops  
polydisperse = droplet size distribution

- Altogether 8 simulations + 1 “single phase jet” simulation with negligible mass loading of small  $d=2\mu m$  tracer particles

# Effect of Mass Loading on the Length of the Potential Core

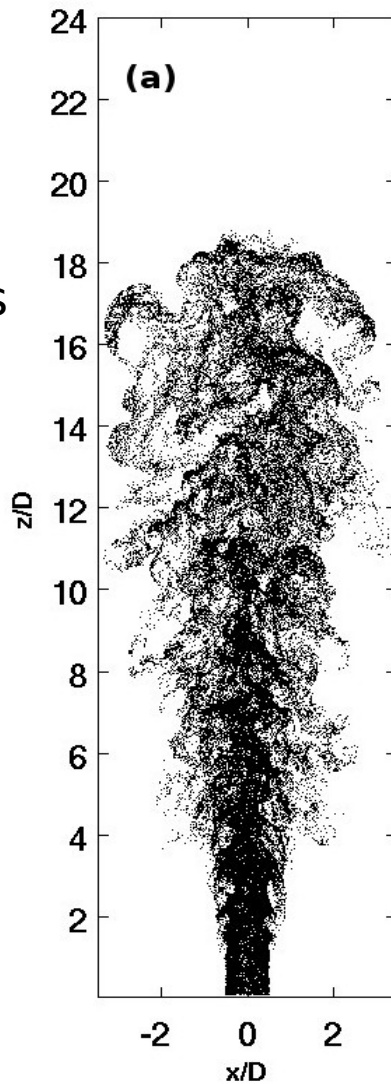


potential  
core: laminar  
interactions



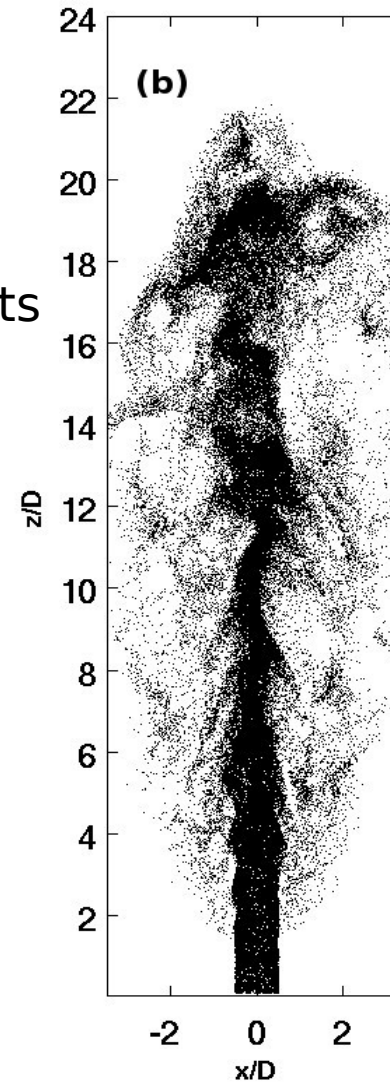
# Different Spray Cloud Shapes

time=1.494ms



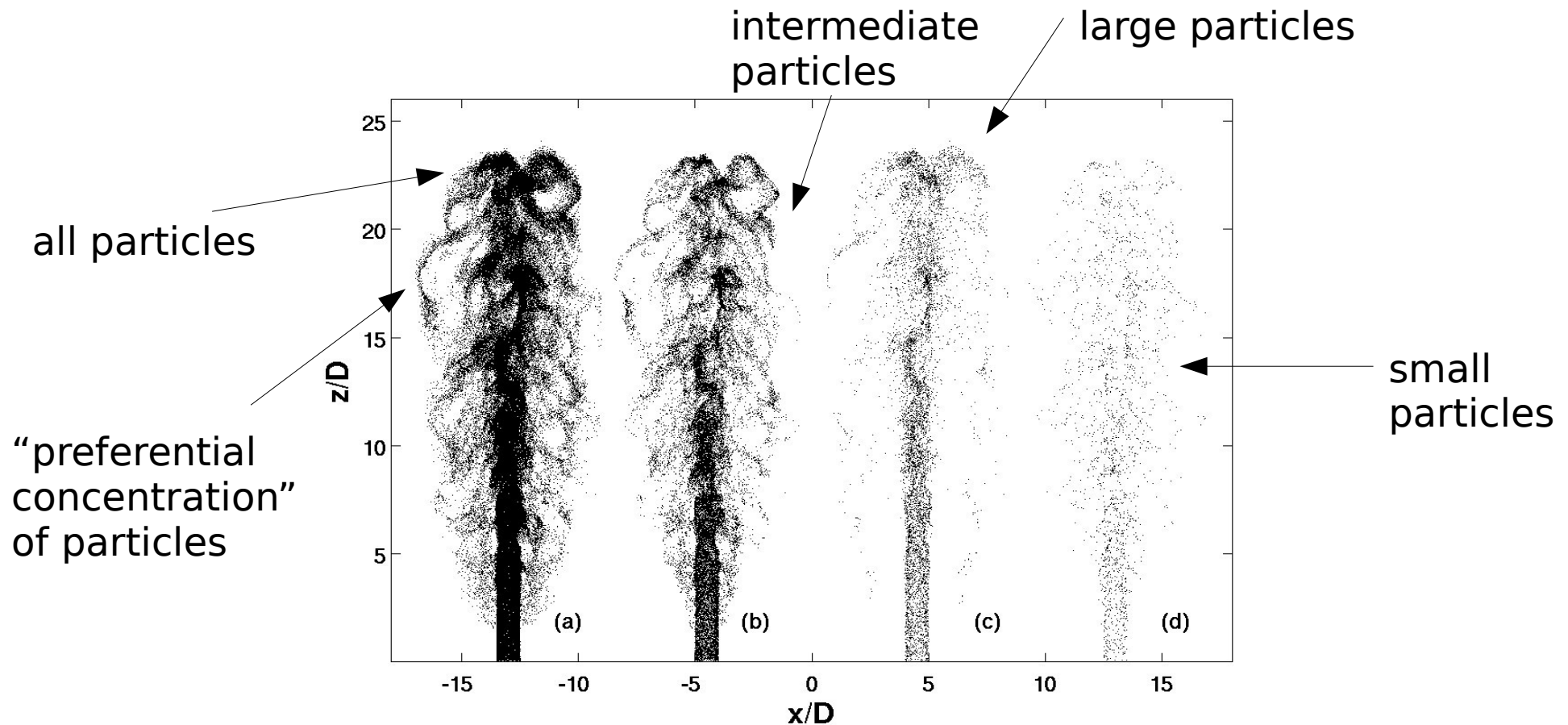
small droplets  
(monodisp.)

large droplets  
(polydisp.)



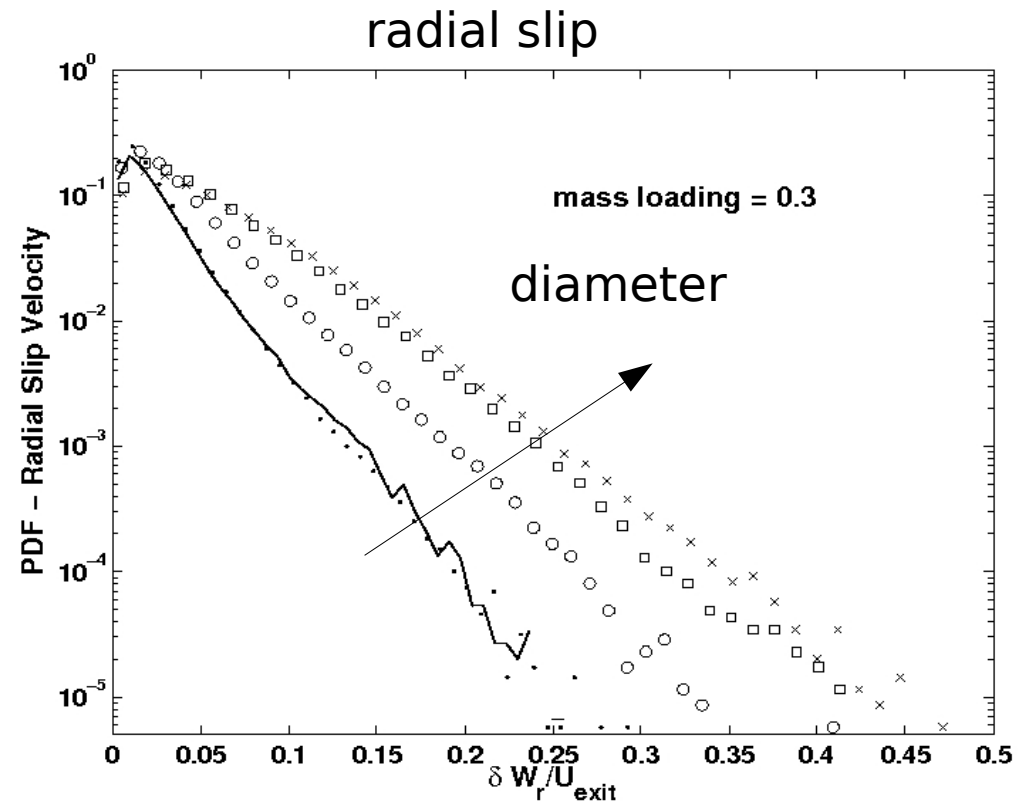
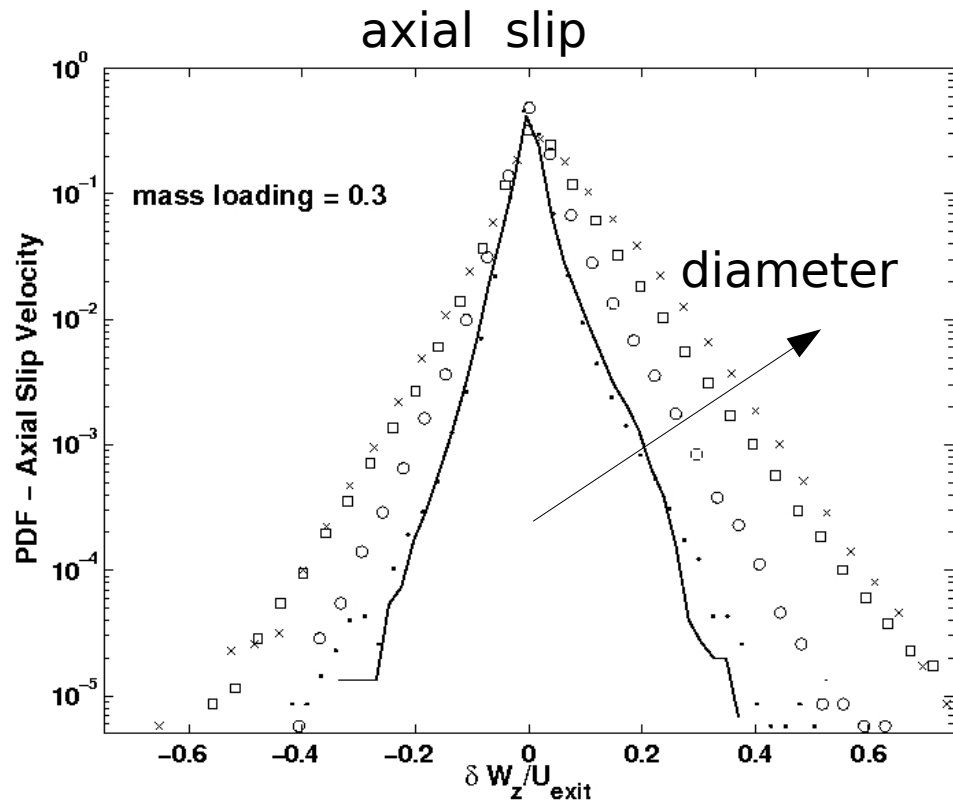
# Superposition of Clouds with Different Stokes Numbers

- Particle Stokes number dominates the trajectory: large Stokes number indicates large inertia and that particle stays close to center

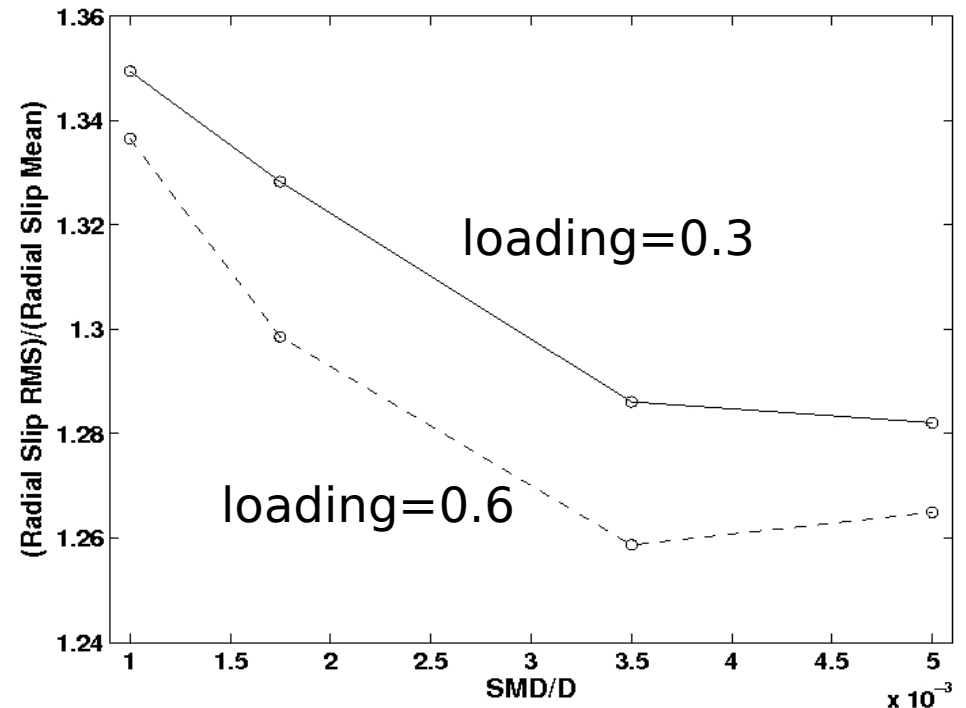
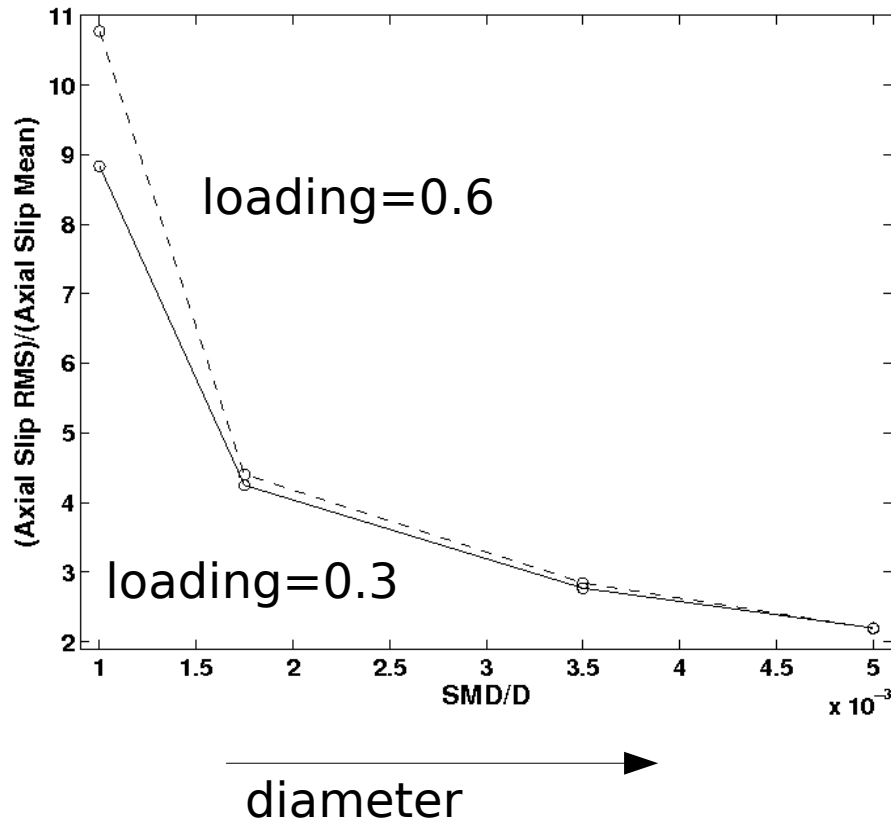


# Probability Density Functions of Slip Velocity

- There is no way the clouds can look different unless the slip velocities  $\delta W$  and their PDFs differ!

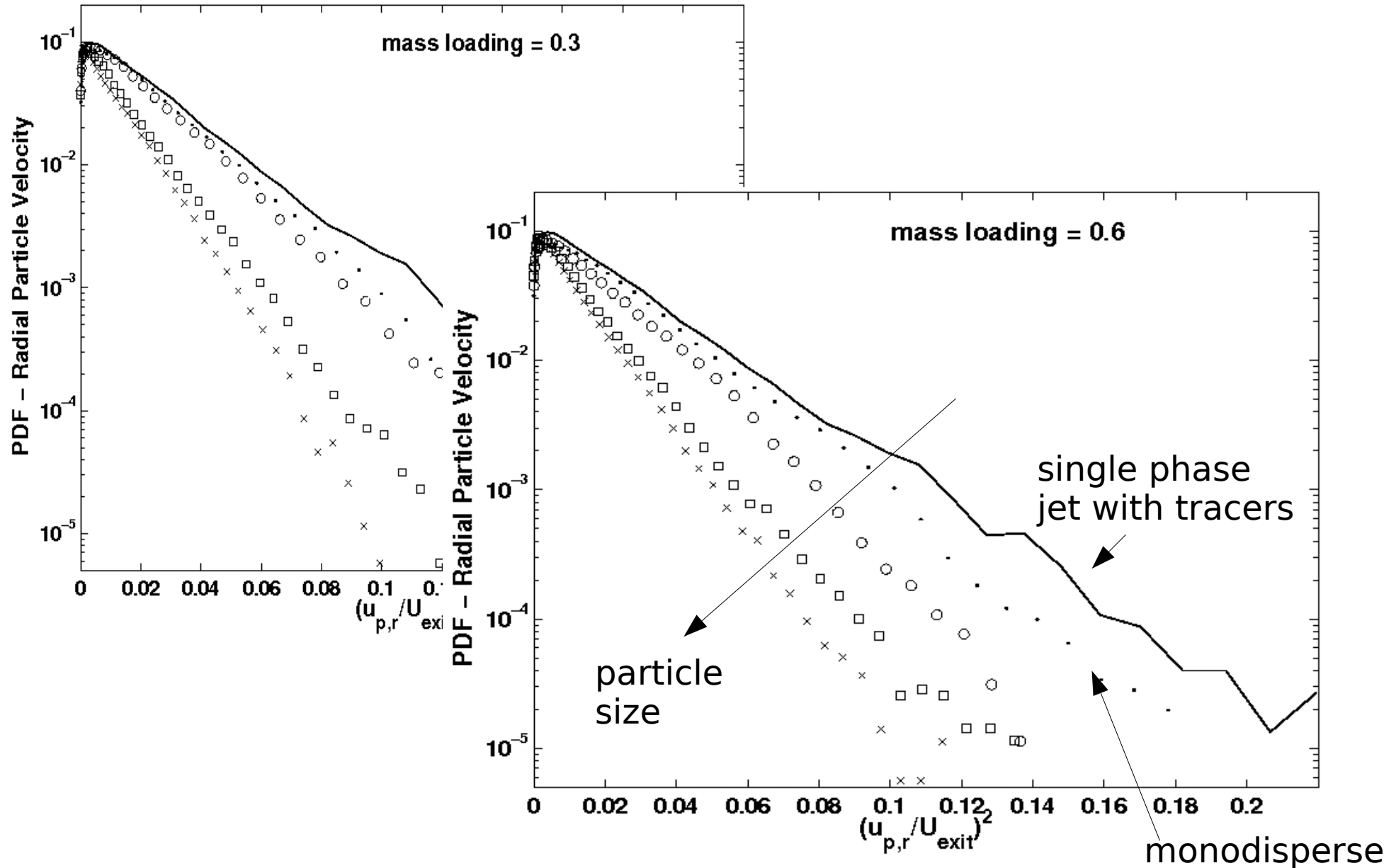


# RMS/MEAN Slip



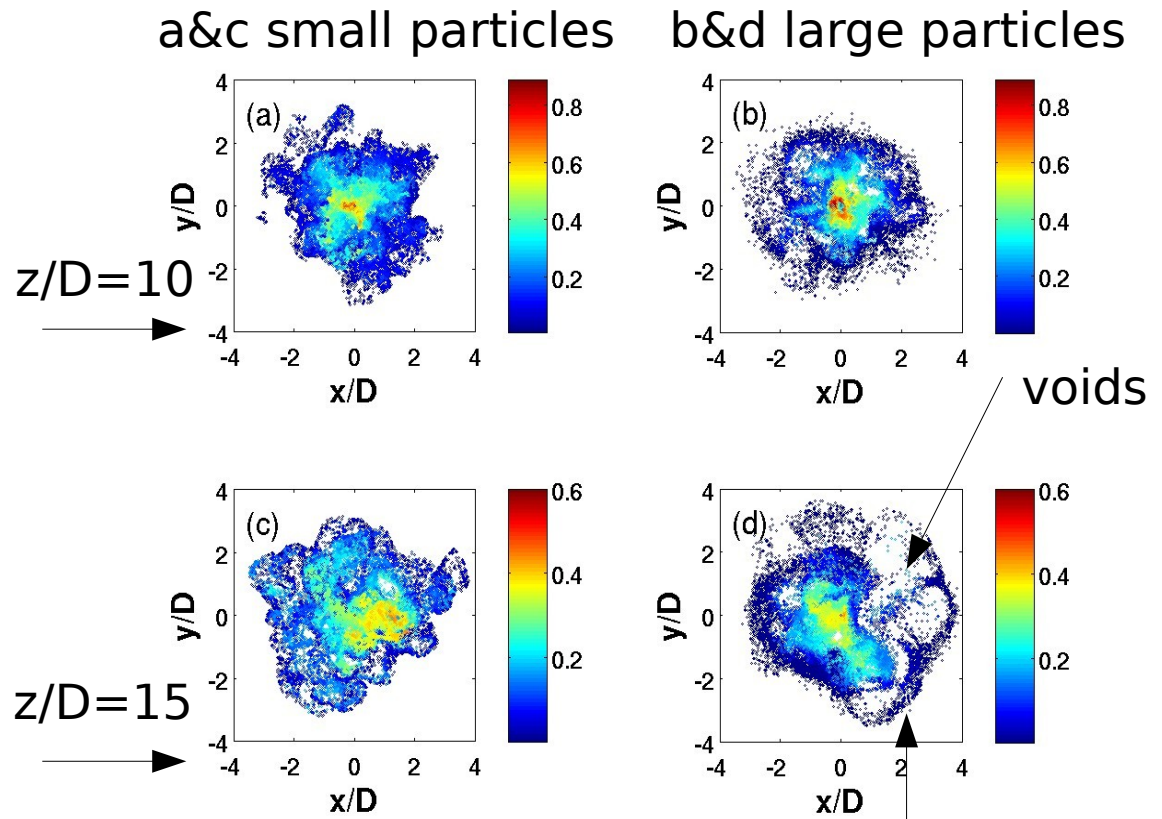
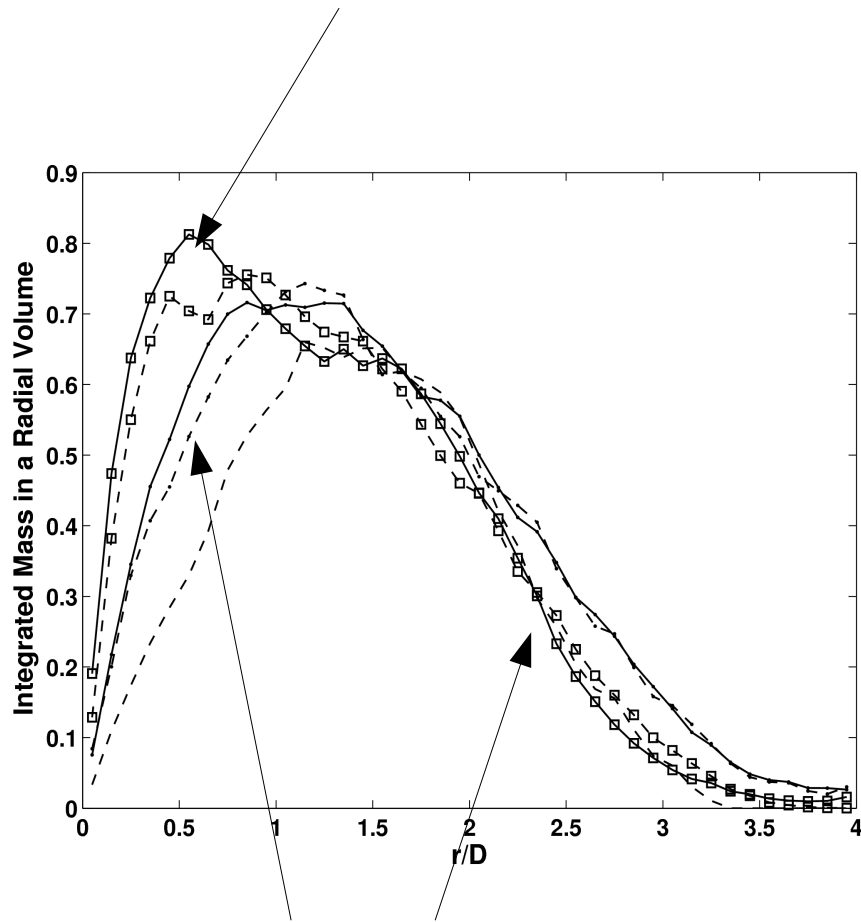
- Small particles have more fluctuating component in the slip velocity: small Stokes number  $\sim$  good temporal response
- Increased loading decreases radial fluctuations and increases axial fluctuations since the axial transport increases

# Probability Density of Radial Particle Velocity



# Mass Transfer and Mixing

large particles  
stay near center



small particles may  
disperse better due to random fluctuations

preferential  
concentration

# Gradient vs Counter Gradient Diffusion

- An often used closure model in turbulence models is the Gradient Diffusion Model (GDM) which associates turbulent fluxes with eddy diffusivity and mean concentration gradient.

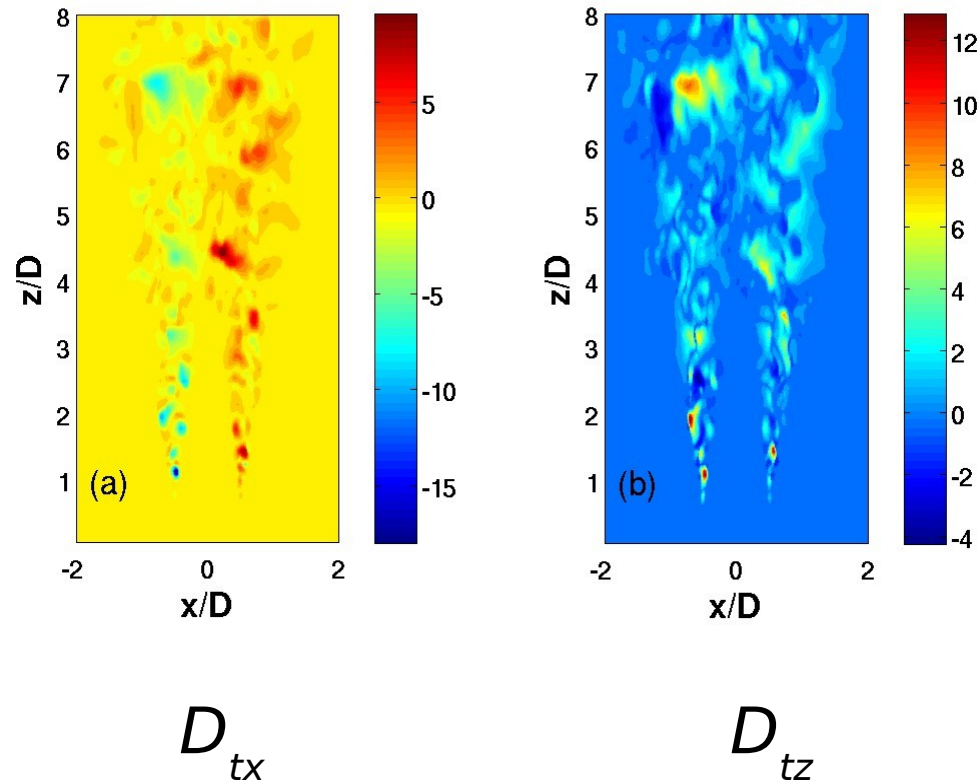
$$\overline{u'_j c'} = -\frac{\nu_t}{Sc_t} \frac{\partial C}{\partial x_j}$$

- The eddy diffusivity can then be solved for:

$$D_{tj} = \frac{-\overline{u'_j c'}}{\frac{\partial C}{\partial x_j}}$$

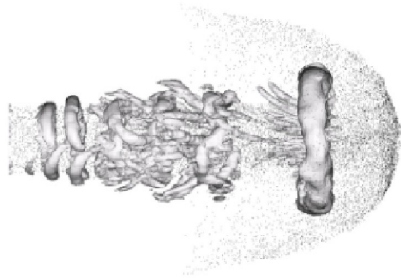
- According to this  $D_{tj}$  should be strictly positive which can be checked for by computation of instantaneous turbulent concentration fluxes.

# Counter Gradient Diffusion is Observed

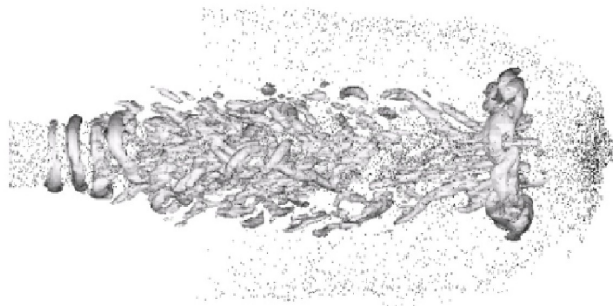


- The eddy diffusivity changes sign from + to -
- Thus the nature of turbulent diffusion is quite different from molecular diffusion and the GDM is not valid in particle laden free shear flows
- Counter gradient diffusion observed for large and small particles and mass loadings

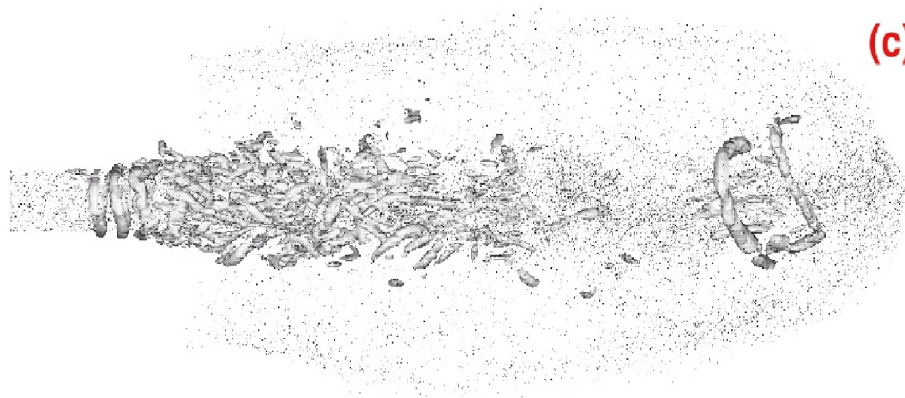
# Look at the Coherent Structures



(a)



(b)



(c)

- Example from Vuorinen et al. (2007): small mass loading of large particles
- Large doughnut shaped tip vortex observed
- Axially oriented vortices in the center
- Breakup of the tip vortex later downstream

Vuorinen, Larmi, Fuchs, *Large-Eddy Simulation of Spray-Originated Turbulence Production and Dissipation*, ICMF-2007, Leipzig, (2007).

# Conclusions

- Spray cloud shape is explained by mean diameter of the droplet distribution and characteristic Stokes numbers – the results are generally in line with several earlier observations on sprays and jets
- PDF's of droplet kinetic energy and slip velocity depend on the mean diameter and mass loading and the tails decay exponentially in the model
- Small scale behavior is reflected to large scale behavior
- Small particles -> follow the fluid motions closely
- Large particles -> particle clustering, high particle concentrations in the center, voids etc...
- Effect of increased mass loading: longer potential cores!!!
- The GDM-assumption could be most problematic in the shear layer
- Large tip vortex and shear layer vorticity are noted and their appearance seems to be sensitive to the droplet size, mass loading and boundary conditions in general.

*THANK YOU FOR YOUR ATTENTION!!!*